

**Radio and Optical Wave Propagation – Prof. Lorenzo Luini**  
**Sample exam**

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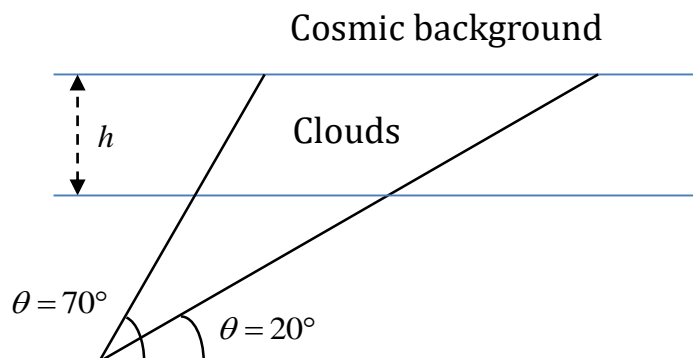
**Exercise 1**

In a cloudy day, a radiometer pointed at  $70^\circ$  elevation measures a brightness temperature  $T_{70} = 110$  K. Calculate:

- 1) The corresponding attenuation, assuming a medium temperature  $T_{mr} = 275$  K.
- 2) The brightness temperature measured and the attenuation estimated by the radiometer in the same day, if pointed at  $20^\circ$  elevation (assume same  $T_{mr}$ ).
- 3) Assuming clouds have a thickness of 1.5 km, compute the equivalent specific attenuation in dB/km.

Consider the cosmic background noise, save for point 3), in which clouds are assumed to be the only source of noise.

**Solution:**



- 1) We know the brightness temperature measured at  $70^\circ$  elevation angle, from which we can derive the attenuation due to clouds ( $T_C = 2.73$  K):

$$T_{70} = T_{mr} (1 - \alpha_{70}) + \alpha_{70} T_C = 110 \text{ K} \rightarrow \alpha_{70} = \frac{T_{70} - T_{mr}}{T_C - T_{mr}} = 0.606 \rightarrow A_{70} = 10 \log_{10} \left( \frac{1}{\alpha_{70}} \right) = 2.18 \text{ dB}$$

2) Assuming horizontal homogeneity for clouds, the attenuation due to clouds along the zenith ( $\theta = 90^\circ$ ) is:

$$A_z = A_{70} \sin 70^\circ = 2.04 \text{ dB}$$

In the same way, at  $20^\circ$ , the attenuation would be:

$$A_{20} = \frac{A_z}{\sin 20^\circ} = 5.96 \text{ dB} \rightarrow \alpha_{20} = 10^{-\frac{A_{20}}{10}} = 0.254 \rightarrow T_{20} = T_{mr} (1 - \alpha_{20}) + \alpha_{20} T_C = 206 \text{ K}$$

3) The zenithal attenuation is given as:

$$A_z = \gamma h$$

Therefore the specific attenuation of clouds  $\gamma$  is:

$$\gamma = \frac{A_z}{h} = 1.36 \text{ dB/km}$$

## Exercise 2

A satellite operating at 12 GHz in RHCP (right hand circular polarization) is seen by the ground station at elevation  $\theta = 20^\circ$ .

Calculate:

- 1) The power transmitted by the satellite  $P_T$  in order counteract the attenuation caused by a slab of uniform rain intensity along the path of 50 mm/h (oblate drops, equi-oriented, minor axis in the vertical direction) and to have an SNR of 5 dB at the ground station.

Assume:

- a) Rain extends up to  $h_R = 3.5$  km in height
- b) Specific attenuation at 12 GHz and  $20^\circ$  elevation is  $\gamma = aR^b$ , where  $a_V = 0.01655$ ,  $b_V = 1.175$  for vertical polarization,  $a_H = 0.01835$ ,  $b_H = 1.219$  for horizontal polarization
- c)  $\gamma = (\gamma_V + \gamma_H) / 2$  for design purposes
- d) On-board antenna  $G_s = 35$  dB
- e) Ground receiver noise temperature  $T_R = 250$  °K
- f) Signal bandwidth  $B = 300$  MHz
- g) Ground antenna gain  $G_g = 50$  dB
- h) Extra attenuation due to gases and clouds  $A_{cg} = 1.5$  dB
- i) Ground-to-satellite distance  $L = 36000$  km
- l) Antennas optimally pointed

- 2) The wave polarization in front of the ground receiving antenna.

### Solution:

- 1) The link budget equation, including the attenuation due to rain, clouds and gases  $A_{dB} = A_{rain} + A_{cg}$ , is:

$$P_R = P_T G_s f_T \left( \frac{\lambda}{4\pi L} \right)^2 G_g f_R A$$

In this equation, both  $f_T$  and  $f_R$  are equal to 1 because the antennas are optimally pointed. We also have information on the required SNR = 5 dB = 3.162, which allows to calculate the minimum received power at the ground station ( $P_N$  is the noise power):

$$SNR = \frac{P_R}{P_N} = \frac{P_R}{T_R B k} \rightarrow P_R = SNR T_R B k = 3.3 \text{ pW}$$

The attenuation due to rain (50 mm/h) can be calculated as:

$$\gamma_H = 2.16 \text{ dB/km}, \gamma_V = 1.64 \text{ dB/km}, \gamma = (\gamma_V + \gamma_H) / 2 = 1.9 \text{ dB/km} \rightarrow A_{rain} = \gamma \frac{h_R}{\sin \theta} = 19.4 \text{ dB}$$

The total attenuation due to rain, clouds and gases is therefore  $A_{dB} = A_{rain} + A_{cg} = 20.9 \text{ dB} \rightarrow$

$$A = 10^{-A_{dB}/10} = 0.0081$$

Thus we can invert the link budget equation to derive  $P_T$ :

$$P_T = \frac{P_R}{G_s \left( \frac{\lambda}{4\pi L} \right)^2 G_g A} = 420.4 \text{ W}$$

where  $\lambda = c/f = 0.025$  m and all the quantities have been converted to linear scale, i.e.  $G_s = 3162.3$ ,  $G_g = 10000$ .

2) The wave gets depolarized by crossing the rain slab, as drops are oblate spheroids. Specifically, the vertical and horizontal components of the original RHCP wave undergo a different attenuation. The RHCP wave is given by the following components:

$$E_V = E_0 \quad \text{and} \quad E_H = -jE_0$$

The two components propagate through the rain slab and at the receiver we have:

$$E_V(L) = E_0 e^{-(\alpha_V + j\beta_V)L_R} = E_0 e^{-\alpha_V L_R} e^{-j\beta_V L_R} = E_0 0.1447 e^{-j\beta L_R}$$

$$E_H(L) = -jE_0 e^{-(\alpha_H + j\beta_H)L_R} = -jE_0 e^{-\alpha_H L_R} e^{-j\beta_H L_R} = -jE_0 0.0784 e^{-j\beta L_R}$$

where the phase constant is the same for both components ( $\beta = 2\pi/\lambda = 251.3 \text{ rad/m}$ ),  $L_R = h_R/\sin(20^\circ) = 10.2 \text{ km}$  is the path travelled through the rain slab and:

$$\alpha_V = \gamma_V \cdot 8.686 \cdot 1000 = 1.889 \cdot 10^{-4} \text{ Np/m}$$

$$\alpha_H = \gamma_H \cdot 8.686 \cdot 1000 = 2.388 \cdot 10^{-4} \text{ Np/m}$$

As is clear, the differential phase between the two components is the same, while they attenuate differently, which leads to a right-end elliptical polarization in front of the receiver: indeed the wave has been depolarized by the rain slab, as the elliptical polarization can be seen as the combination of a RHCP and LHCP waves.

### Exercise 3

Calculate the additional attenuation of a wave travelling through the ionosphere for a zenithal pointing, assuming that the collision rate is  $\nu = 10^4$  collisions/s and that the electrons are present and with a constant content  $N$  from 100 km up to 400 km. The critical frequency is 900 kHz and the wave frequency is 10 MHz.

#### Solution:

From the critical frequency, we can derive the electron content  $N$ :

$$f_c \approx 9\sqrt{N} \Rightarrow N = \frac{f_c^2}{81} \approx 10^{10} \text{ e/m}^3$$

The equivalent conductivity of the ionosphere is:

$$\sigma = \frac{Ne^2\nu}{m(\nu^2 + \omega^2)} = 7.2 \cdot 10^{-10} \text{ S/m}$$

where  $m = 9 \cdot 10^{-31}$  kg is the mass of the electron and  $e = -1.6 \cdot 10^{-19}$  C is its charge.

The plasma angular frequency (squared) is:

$$\omega_p^2 = \frac{Ne^2}{m\epsilon_0} = 3.2 \cdot 10^{13} \text{ rad}^2/\text{m}^2$$

from which we can calculate the equivalent relative permittivity of the ionosphere:

$$\epsilon_r = 1 - \frac{\omega_p^2}{\nu^2 + \omega^2} = 0.9919$$

The propagation constant thus is:

$$\gamma = \alpha + j\beta = \sqrt{j\omega\mu(\sigma + j\omega\epsilon)} = 1.36 \cdot 10^{-7} + j0.21 \text{ 1/m}$$

The total path attenuation is:

$$\alpha_{dB} = \alpha \cdot 8.686 \cdot 1000 = 0.0012 \text{ dB/km}$$

$$A = \alpha_{dB} L = \alpha_{dB} (400 - 100) = 0.355 \text{ dB}$$

#### Exercise 4

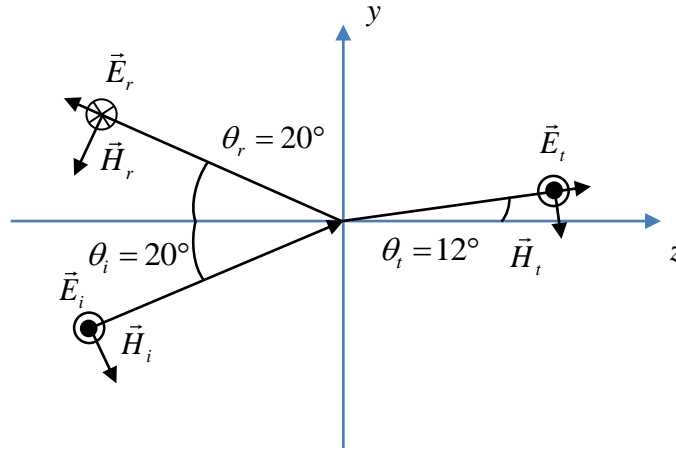
A uniform plane wave at  $f = 20$  GHz with linear polarization propagates from a medium with  $\epsilon_{r1} = 2$  to another medium with different electromagnetic features. The electric field of the incident wave, whose value is  $E_0 = 10$  V/m, is parallel to the boundary between two media and the incident angle is  $20^\circ$ .

Calculate:

- 1) The electric field of the incident field
- 2) The density power carried by the wave travelling in the second medium with a refraction angle of  $12^\circ$
- 3) The magnetic field of the reflected wave

Assume both media are lossless and non-magnetic.

**Solution:**



1) Knowing the incidence and refraction angles and exploiting Snells' law, we can derive the relative electric permittivity of the second medium:

$$\sqrt{\epsilon_{r1}} \sin \theta_i \Rightarrow \sqrt{\epsilon_{r2}} \sin \theta_t \rightarrow \epsilon_{r2} = \left( \frac{\sqrt{\epsilon_{r1}} \sin \theta_i}{\sin \theta_t} \right)^2 = 5.4$$

Therefore the expression for the incident electric field will be:

$$\vec{E}_i = (-10\vec{\mu}_x) e^{-j\beta_1(\cos \theta_i \vec{\mu}_z + \sin \theta_i \vec{\mu}_y)}$$

$$\text{where } \beta_1 = \frac{2\pi f}{c} \sqrt{\epsilon_{r1}} = 592.4 \text{ rad/m}$$

2) To calculate the power density flowing in the second medium, we first need to calculate the reflection coefficient (TE case) and the intrinsic impedance of the two media:

$$\eta_1^{TE} = \frac{\eta_0}{\sqrt{\epsilon_{r1}}} \frac{1}{\cos \theta_i} = 283.7 \text{ } \Omega$$

$$\eta_2^{TE} = \frac{\eta_0}{\sqrt{\epsilon_{r2}}} \frac{1}{\cos \theta_t} = 165.9 \text{ } \Omega$$

$$\Gamma^{TE} = \frac{\eta_2^{TE} - \eta_1^{TE}}{\eta_2^{TE} + \eta_1^{TE}} = -0.26$$

We can now calculate the power density flowing into medium as:

$$S_1 = \frac{1}{2} \frac{|\vec{E}_i|^2}{\eta_0 / \sqrt{2}} = 0.188 \text{ W/m}^2$$

$$S_2^z = S_1^z (1 - |\Gamma^{TE}|^2) \Rightarrow S_2 = \frac{S_1 \cos \theta_i (1 - |\Gamma^{TE}|^2)}{\cos \theta_t} = 0.168 \text{ W/m}^2$$

3) For the reflected magnetic field, we first need to calculate the reflected electric field as:

$$\vec{E}_r = \vec{E}_i \Gamma^{TE} = 2.6 \vec{\mu}_x \text{ V/m}$$

The absolute value of the reflected magnetic field will be:

$$|\vec{H}_r| = |\vec{E}_r| / \eta_1 = 0.0098 \text{ A/m}$$

Finally, the full expression for the magnetic field will be:

$$\vec{H}_r = \left( -|\vec{H}_r| \cos \theta_i \vec{\mu}_y - |\vec{H}_r| \sin \theta_i \vec{\mu}_z \right) e^{-j\beta(-\cos \theta_r z + \sin \theta_r y)} \text{ A/m}$$

## Exercise 5

Consider a transmitting antenna installed on top of a pole whose height is  $h = 70$  m, with an output power  $P_T = 5$  W (antenna gain  $G_{tx} = 40$  dB).

Calculate:

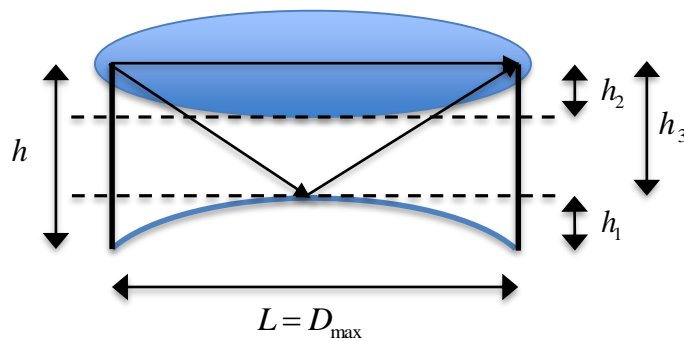
- 1) The maximum distance  $D_{max}$  covered by the antenna in standard atmospheric conditions.
- 2) Assuming now a receiver at distance  $D_{max}$  from the transmitter, at the same height  $h = 70$  m (antenna with the same gain  $G_{tx} = 40$  dB), calculate the frequency of the link by knowing that, in the propagation conditions at point 1), the distance ground-direct ray is 2.5 times the semi-minor axis of the first Fresnel's ellipsoid.
- 3) The power received considering only the direct ray (assume optimal antenna pointing).
- 4) Considering also the reflected ray (assume perfectly reflecting surface), is the combination of the two rays at the receiver constructive or destructive?

### Solution:

1) The maximum distance covered can be calculated as (standard atmosphere  $\rightarrow$  equivalent Earth radius factor = 4/3):

$$D_{max} = \sqrt{2hR_{eq}} = \sqrt{2hkR_E} = \sqrt{2h \frac{4}{3} R_E} = 34.4 \text{ km}$$

2)



From the text we know that  $h_3 = 2.5 h_2$ , being  $h_2$  the semi-minor axis of the first Fresnel's ellipsoid, i.e.  $h_2 = \sqrt{\lambda L}/2$ . At  $L = D_{max} = 34.3$  km, we have:

$$h_1 = \frac{1}{2} \frac{(L/2)^2}{R_{eq}} = 17.4 \text{ m}$$

$$h_3 = 2.5h_2 = h - h_1 = 52.6 \text{ m} \rightarrow h_2 = \frac{h_3}{2.5} = 21 \text{ m} \rightarrow \lambda = \frac{(2h_2)^2}{L} = 0.0513 \text{ m}$$

Therefore:

$$f = \frac{c}{\lambda} = 5.85 \text{ GHz}$$

3) Using the link budget:

$$P_R = P_T G_{tx} \left( \frac{\lambda}{4\pi L} \right)^2 G_{rx} = 7 \mu\text{W}$$

4) Defining  $E_0$  as the electric field at the receiver and associated to the direct ray, the combination of the direct and reflected rays can be assessed as:

$$E = E_0(1 + e^{-j\beta\delta}) \text{ V/m}$$



where  $\delta$  is the differential path, i.e. the difference between the path travelled along the reflected ray and the one travelled along the direct ray. This quantity can be well approximated by:

$$\delta = \frac{2h_3^2}{L} = 0.1609 \text{ m}$$

Therefore:

$$E = E_0(1 + e^{-j\beta\delta}) = E_0(1.6531 - j0.7573) \text{ V/m} \rightarrow |E| = |E_0|1.818 \text{ V/m}$$

The combination of the two rays is therefore constructive.