

**Telecommunication Systems – Prof. L. Luini,  
January 29<sup>th</sup>, 2021**

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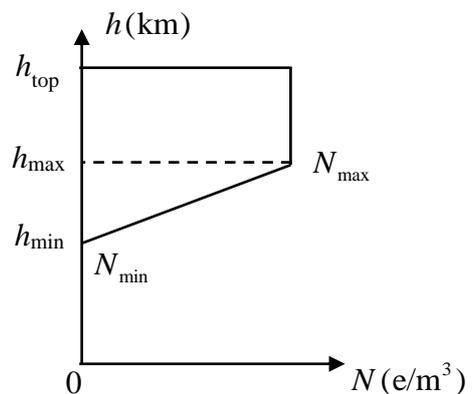
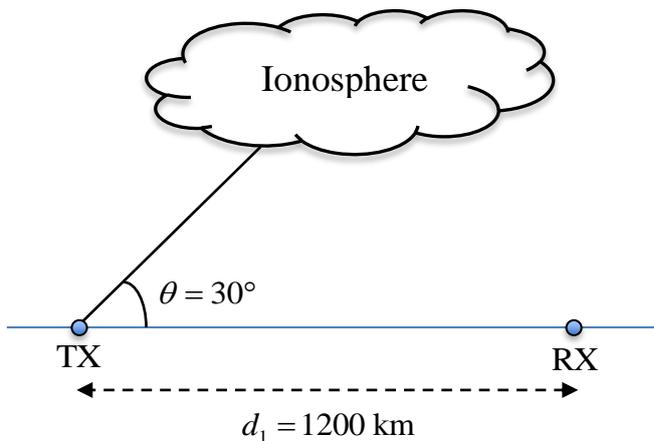
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**Problem 1**

Making reference to the figure below, the transmitter TX needs to reach the user RX, which is at distance  $d_1 = 1200$  km, by exploiting the effects induced by the ionosphere. The elevation angle of TX is  $\theta = 30^\circ$ . The ionosphere can be modelled with the electron density profile sketched in the figure (right side), where  $N_{\min} = 4 \times 10^2 \text{ e/m}^3$ ,  $N_{\max} = 4 \times 10^{12} \text{ e/m}^3$ ,  $h_{\min} = 100$  km,  $h_{\max} = 400$  km and  $h_{\text{top}} = 700$  km.

- 1) Calculate the link frequency  $f_1$  to reach RX.
- 2) Assuming now that RX can move further from TX (keeping the same elevation angle), calculate the maximum distance  $d_2$  at which RX can be reached, as well as the corresponding link frequency  $f_2$ .
- 3) What happens if the link frequency increases beyond  $f_2$ ?

Assume: the virtual reflection height  $h_V$  is 1.2 of  $h_R$ , the height at which the wave is actually reflected.



**Solution:**

1) Considering the figure below, given the distance between the TX and RX and the elevation angle, the virtual reflection height  $h_V$  is given by:

$$h_{V1} = \frac{d_1}{2} \tan \theta = 346.4 \text{ km}$$

The actual reflection occurs at  $h_{R1} = h_{V1}/1.2 = 288.7 \text{ km}$ . The value of electron density  $N$  associated to  $h_{R1}$  can be obtained from the following linear expression (easily derivable from the profile sketch):

$$N = \frac{N_{\max} - N_{\min}}{h_{\max} - h_{\min}} (h - h_{\min}) + N_{\min}$$

Thus, for  $h = h_{R1}$ , we obtain:

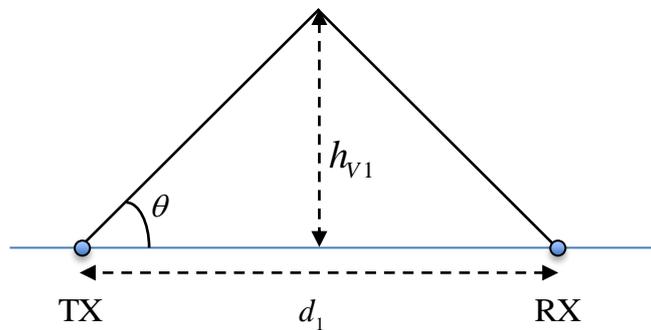
$$N_1 = 2.5 \times 10^{12} \text{ e/m}^3$$

The elevation angle and the frequency are linked by the following expression:

$$\cos \theta = \sqrt{1 - \left( \frac{f_p}{f_1} \right)^2} = \sqrt{1 - \left( \frac{9\sqrt{N_1}}{f_1} \right)^2}$$

Solving for the frequency  $f_1$ , we obtain:

$$f_1 = \frac{\sqrt{81N_1}}{\sqrt{1 - [\cos(\theta)]^2}} \approx 28.6 \text{ MHz}$$



2) The maximum distance  $d_2$  between TX and RX corresponds to the point of maximum reflection along the profile, which corresponds to  $h_{\max}$ , i.e. the height where the maximum electron density  $N_{\max}$  is found. The maximum distance can be easily calculated as:

$$d_2 = \frac{2h_{V2}}{\tan \theta} = \frac{2 \cdot 1.2 \cdot h_{\max}}{\tan \theta} = 1662.3 \text{ km}$$

The correspondent link frequency needed to achieve reflection at such height is:

$$f_2 = \frac{\sqrt{81N_{\max}}}{\sqrt{1 - [\cos(\theta)]^2}} \approx 36 \text{ MHz}$$

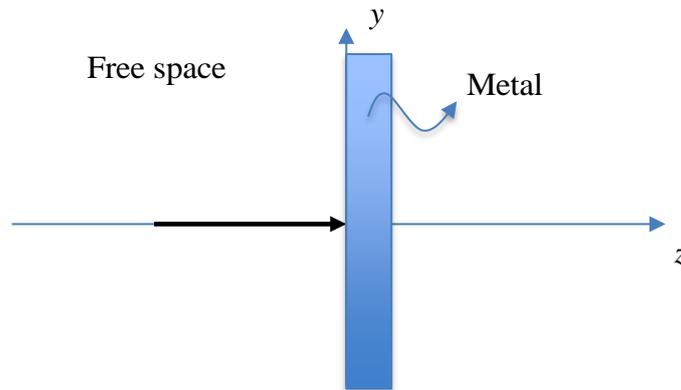
3) If the operational frequency exceeds  $f_2$ , the wave crosses the ionosphere.

## Problem 2

A plane sinusoidal EM wave propagates from vacuum into a perfect metallic surface (perfect electric conductor) with orthogonal incidence (assume both  $\epsilon_r = 1$  and  $\mu_r = 1$  for this material). The expression for the electric field is:

$$\vec{E} = (j\vec{\mu}_y + \vec{\mu}_x)e^{-j188.62z} \text{ V/m}$$

- 1) Determine the frequency of the EM wave.
- 2) Determine the polarization of the incident EM wave.
- 3) Determine the polarization of the reflected wave.



### Solution:

1) The frequency of the incident EM wave can be derived from the phase constant  $\beta = 188.62$  rad/m:

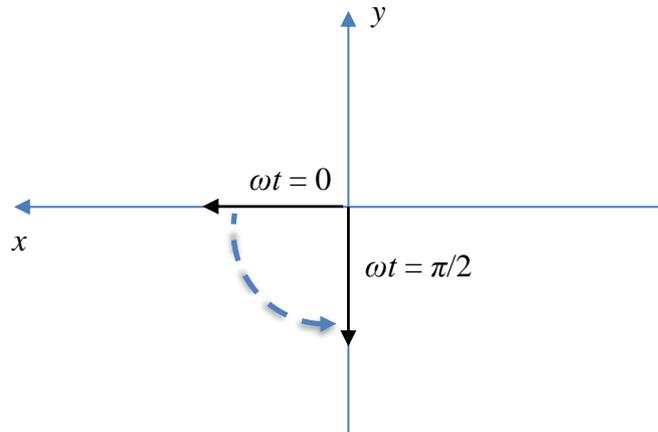
$$\beta = \omega\sqrt{\mu_0\epsilon_0} = \frac{2\pi f}{c} \Rightarrow f = \frac{c\beta}{2\pi} = 9 \text{ GHz}$$

2) The polarization of the incident wave is circular because the two TE and TM components have the same amplitude (1 V/m) and a phase shift of  $\pi/2$ . Setting  $z$  to 0, and expressing the dependence on time, we can easily understand the electric field rotation direction:

$$\vec{E}(t) = \text{Re}\left\{\left[j\vec{\mu}_y + \vec{\mu}_x\right]e^{j\omega t}\right\} = \cos\left(\omega t + \frac{\pi}{2}\right)\vec{\mu}_y + \cos(\omega t)\vec{\mu}_x \text{ V/m}$$

Thus, for  $t = 0 \rightarrow \vec{E}\Big|_{\omega t=0} = \vec{\mu}_x \text{ V/m}$

Thus, for  $\omega t = \pi/2 \rightarrow \vec{E}\Big|_{\omega t=\pi/2} = -\vec{\mu}_y \text{ V/m}$



The polarization of the incident wave is LHCP.

3) When a plane wave hits a perfect electric conductor (PEC), the reflection coefficient is always  $\Gamma = -1$ . Therefore, both components of the electric field will change their sign, and the wave plane direction will also change. This is shown in the equation of the reflected wave:

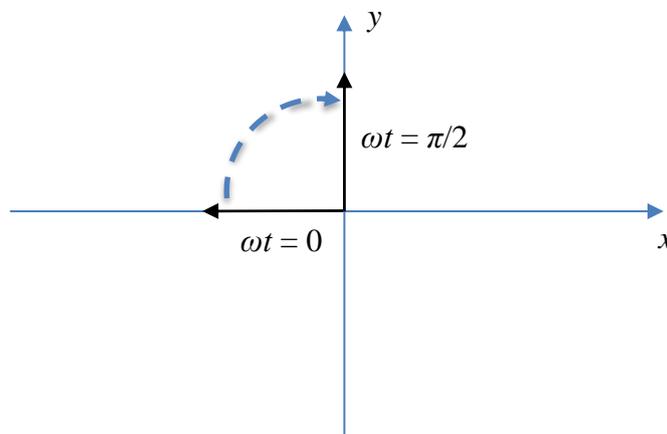
$$\vec{E}_R = (-j\vec{\mu}_y - \vec{\mu}_x) e^{j188.62z} \text{ V/m}$$

The polarization of the reflected wave is still circular; also, repeating the same reasoning above, we obtain:

$$\vec{E}_R(t) = \text{Re} \left\{ [-j\vec{\mu}_y - \vec{\mu}_x] e^{j\omega t} \right\} = -\cos\left(\omega t + \frac{\pi}{2}\right) \vec{\mu}_y - \cos(\omega t) \vec{\mu}_x \text{ V/m}$$

Thus, for  $t = 0 \rightarrow \vec{E}_R|_{\omega t=0} = -\vec{\mu}_x \text{ V/m}$

Thus, for  $\omega t = \pi/2 \rightarrow \vec{E}_R|_{\omega t=\pi/2} = \vec{\mu}_y \text{ V/m}$



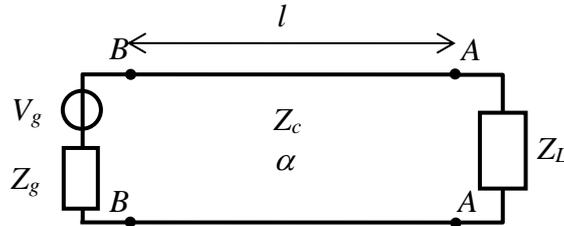
The polarization of the incident wave is RHCP. Note that, as the direction of the reflected plane wave changes (direction =  $-z$ ), so does also the view of the axes, with the  $x$ -axis pointing to the right side this time.

### Problem 3

A transmitter with voltage  $V_g = 100$  V (sinusoidal regime) and internal impedance  $Z_g = 50 \Omega$  is connected to a load  $Z_L = -j10 \Omega$  by a transmission line with characteristic impedance  $Z_C = 100 \Omega$  and attenuation coefficient  $\alpha = 30$  dB/km . The line length is  $l = 30$  m and the frequency is  $f = 300$  MHz.

Calculate:

- 1) The power absorbed by the load.
- 2) The value of  $Z_L$  to maximize  $P_L$ , the power transferred to the load.
- 3) The value of  $P_L$  and the power absorbed by the line, for the conditions at point 2).



### Solution

1) As the load is imaginary (it corresponds to a capacitor), no power will be absorbed by  $Z_L$ .

2) The value of  $Z_L$  maximizing the power absorbed by the load is:

$$Z_L = 100 \Omega = Z_C$$

In fact, in this case, the load will guarantee at least a partial match of the load with the transmission line.

3) In the conditions at point 2), the power absorbed by the load is given by:

$$P_L = P_{AV} (1 - |\Gamma_g|^2) e^{-2\alpha l} \text{ W}$$

$$P_{AV} = \frac{|V_g|^2}{8 \operatorname{Re}\{Z_g\}} = 25 \text{ W}$$

The attenuation coefficient must be first converted to Np/m:

$$\alpha = \frac{30}{8.686 \cdot 1000} = 0.0035 \text{ Np/m}$$

Finally:

$$\Gamma_g = \frac{Z_{BB} - Z_g}{Z_{BB} + Z_g} = \frac{Z_L - Z_g}{Z_L + Z_g} = 0.334$$

Note that in this case  $Z_{BB} = Z_L$  as the load is matched to the line.

Therefore:

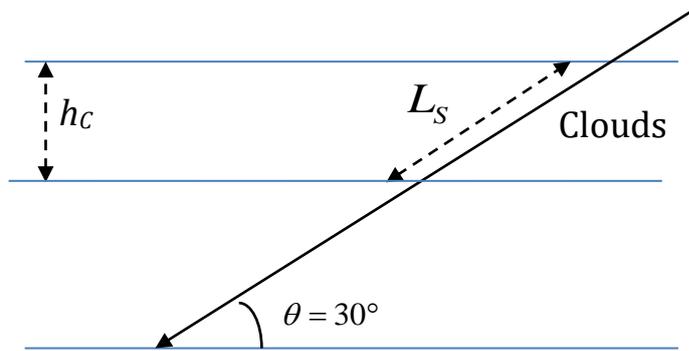
$$P_L = 18.1 \text{ W}$$

Finally, the power absorbed by the line is:

$$P_{line} = P_{AV} (1 - |\Gamma_g|^2) - P_L = 4.1 \text{ W}$$

#### Problem 4

Consider a link from a GEO satellite to a ground station (elevation angle  $\theta = 30^\circ$ ) and assume that, initially, there is no atmospheric attenuation: associate this condition to the reference signal-to-noise ratio  $SNR_0$ . Evaluate the decrease in the SNR (as a function of  $SNR_0$ ) when clouds begin to affect the link. Consider: no cosmic background noise, link frequency  $f = 50$  GHz, cloud thickness  $h_C = 2$  km,  $T_{mr}$  for clouds =  $0^\circ\text{C}$ , specific cloud attenuation (constant vertically and horizontally)  $\gamma = 1.46$  dB/km, receiver internal noise temperature  $T_R = 300$  K.



#### Solution:

The SNR is given by (no attenuation case):

$$SNR_0 = \frac{P_R}{P_N} = \frac{P_T G_T f_T (\lambda/4\pi D)^2 G_R f_R}{k T_R B}$$

where  $D$  is the distance between the satellite and the ground station and  $B$  is the RX bandwidth.

With cloud attenuation, the SNR changes to:

$$SNR_1 = \frac{P_R A_C}{P_N^C} = \frac{P_T G_T f_T (\lambda/4\pi D)^2 G_R f_R A_C}{k T_C B}$$

where  $A_C$  is the rain induced attenuation and  $T_C$  is the RX noise temperature in case of clouds.

As a matter of fact, the SNR decreases both because of the additional attenuation induced by clouds and because of the increase in the noise received by the antenna.

As for  $A_C$ :

$$(A_C)_{dB} = \gamma L_S = \gamma \frac{h_C}{\sin \theta} \approx 5.84 \text{ dB} \rightarrow A_C = 10^{-(5.84/10)} = 0.26$$

As for  $T_C$ :

$$T_C = T_R + T_A = T_R + (1 - A_C) T_{mr} \approx 502 \text{ K}$$

In cloudy conditions, the RX noise increases by a factor 1.67 if compared to the initial conditions; i.e.:

$$\frac{T_C}{T_R} = \frac{T_R + (1 - A_C)T_{mr}}{T_R} \approx 1.67 \rightarrow T_C = 1.67T_R$$

As a result:

$$SNR_1 = \frac{P_R A_C}{P_N^C} = \frac{A_C}{1.67} \frac{P_T G_T f_T (\lambda/4\pi D)^2 G_R f_R}{k T_R B} = \frac{A_C}{1.67} SNR_0 = 0.16 SNR_0$$